

## Lecture 7 *characteristic classes*

In the previous lectures, we have seen cases where fiber bundles are characterized by integers. For example, monopole bundles on  $S^2$  are classified by the monopole number  $n$  which tells us how the  $U(1)$  fibers over the upper hemisphere and the lower hemisphere are glued together. Generally speaking, for a vector bundle on a manifold  $M$ , a characteristic class associates a cohomology class of  $M$ .

### Invariant Polynomials

Characteristic classes are constructed as polynomials of the curvature  $F = dA + A \wedge A$ . Under gauge transformation,  $F$  transforms as  $F \rightarrow \Omega^{-1} F \Omega$ , where  $\Omega$  is a map from the manifold  $M$  to the gauge group (structure group)  $G$ . In the following, we consider the cases where  $G = U(k)$  and  $SO(2r)$ . To construct characteristic classes, we need to introduce invariant polynomials of matrixes. We look for a function  $P(X)$  of a matrix  $X$  that is invariant under the conjugation,  $P(\Omega^{-1} X \Omega) = P(X)$ . We consider two cases:

- (1)  $X$  is a  $k \times k$  hermitial matrix and  $\Omega \in U(k)$ . This will be used when  $E$  is a complex vector bundle.
- (2)  $X$  being a  $2r \times 2r$  real anti-symmetric matrix and  $\Omega \in SO(2r)$ . This will be used when  $E$  is a real vector bundle.

Examples of invariant polynomials are  $\text{tr} X^m$  ( $m = 1, 2, \dots$ ) and  $\det X$ . In fact we can use these to construct a nice basis. The following two are particularly useful (we are using the notation in the case of  $G = U(k)$ ):

- (1)  $\sigma_i(X)$  defined by

$$\det(1 + tX) = 1 + t\sigma_1(X) + t^2\sigma_2(X) + \dots + t^k\sigma_k(X).$$

- (2)  $s_i(X)$  defined by

$$s_i(X) = \text{tr} X^i, \quad (i = 1, \dots, k).$$

They are related to each other by Newton's formula,

$$s_1 = \sigma_1, \quad s_2 = \sigma_1^2 - \sigma_2, \quad s_3 = \sigma_1^3 - 3\sigma_1\sigma_2 + 3\sigma_3, \dots$$

We can also express the invariant polynomials in terms of eigenvalues. If  $X$  is a hermitian matrix, we can diagonalize it with eigenvalues  $x_1, \dots, x_k$ . Then,

$$\prod_{i=1}^k (1 + tx_i) = 1 + t\sigma_1(x) + \dots + t^k\sigma_k(x).$$

Similarly,

$$s_i(x) = \sum_{j=1}^k (x_j)^i.$$

### Characteristic Classes

If  $P_i(X)$  is an invariant polynomial of degree  $i$ , we can use the curvature 2-form  $F$  to define a  $2i$ -form  $P_i(F)$ . By construction, it is invariant under the gauge transformation,  $F \rightarrow \Omega^{-1}F\Omega$ . It is also a closed form. To see this, we first note that  $F$  satisfies the Bianchi identity,

$$dF + [A, F] = 0.$$

Question 1: Verify the above identity. To show this, one has to use the Jacobi identity

$$[a, [b, c]] + \text{cyclic perm.} = 0.$$

Using this, we find

$$\begin{aligned} d\text{tr}F^i &= \text{tr} \{dFF^{i-1} + FdFF^{i-1} + \dots + F^{i-1}dF\} \\ &= \text{tr} \{(dF + [A, F])F^{i-1} + \dots + F^{i-1}(dF + [A, F])\} \\ &= 0 \end{aligned} \tag{1}$$

$$\tag{2}$$

Thus,  $P_i(F)$  is potentially a non-trivial element of  $H^{2i}(M)$ .

Moreover,  $P_i(F)$  is invariant under continuous deformation of the gauge field  $A$  as an element of  $H^{2i}(M)$ . Suppose we change  $A \rightarrow A + \eta$  with  $\eta$  being an infinitesimal one-form. Note that, although  $A$  transforms inhomogeneously as  $A \rightarrow \Omega^{-1}A\Omega + \Omega^{-1}d\Omega$ , the one-form transforms homogeneously,  $\eta \rightarrow \Omega^{-1}\eta\Omega$ .

Under this deformation,  $F$  changes by  $\delta F = d\eta + [A, \eta]$ . Therefore,

$$\begin{aligned} \delta\text{tr}F^i &= \text{tr} \{(d\eta + [A, \eta])F^{i-1} + \dots + F^{i-1}(d\eta + [A, \eta])\} \\ &= i\text{tr} \{(d\eta + [A, \eta])F^{i-1}\} \\ &= i\text{tr} \{d\eta F^{i-1} + \eta dFF^{i-2} \dots + \eta F^{i-2}dF\} \\ &= id\text{tr}(\eta F^{i-1}). \end{aligned} \tag{3}$$

$$\tag{4}$$

Since both  $\eta$  and  $F$  transform homogeneously under the gauge transformation,  $\text{tr}(\eta F^{i-1})$  is a well-defined  $(2i-1)$ -form. Thus, under any infinitesimal deformation,  $P_i(F)$  changes by an exact form. Thus,  $P_i(F)$  depends only on the type of the bundle  $E$  and not on a specific type of the connection  $A$  on  $E$ . For this reason, we sometime write a characteristic class as a function of  $E$ .

### Chern classes, Chern characters, etc

Among characteristic classes for an  $U(n)$  bundle, we have the Chern classes and the Chern characters.

The **Chern classes**,  $c_i \in H^{2i}(M)$  ( $i = 0, 1, 2, \dots, k$ ), are defined by

$$\det \left( 1 + \frac{\sqrt{-1}}{2\pi} F \right) = c_0 + c_1 + c_2 + \dots$$

For example,

$$c_0 = 1, \quad c_1 = \frac{\sqrt{-1}}{2\pi} \text{tr}F, \quad c_2 = -\frac{1}{8\pi^2} (\text{tr}F \wedge \text{tr}F - \text{tr}F \wedge F), \quad \dots$$

If the holonomy is in  $SU(k) \in U(k)$ , we have a trivial first Chern class  $c_1 = 0$ .

The sum  $c = c_0 + c_1 + c_2 + \dots$  is called the total Chern class. One of the important properties of the Chern classes is that it behaves nicely when we take a direct sum of vector bundles  $E_1, E_2 \rightarrow E_1 \oplus E_2$  as,

$$c(E_1 \oplus E_2) = c(E_1) \wedge c(E_2).$$

On the other hand, it does not behave nicely under the direct product  $E_1 \otimes E_2$ .

The **Chern characters**,  $ch_i(F) \in H^{2i}(M)$  ( $i = 0, 1, 2, \dots$ ), are defined by

$$ch_i(F) = \frac{1}{i!} \text{tr} \left( \frac{\sqrt{-1}F}{2\pi} \right)^i.$$

We can also write it as

$$ch(F) = ch_0 + ch_1 + \dots = \text{tr} \exp \left( \frac{\sqrt{-1}F}{2\pi} \right).$$

The Chern characters behave nicely under both the direct sum and direct product as,

$$\begin{aligned} ch(E_1 \oplus E_2) &= ch(E_1) + ch(E_2), \\ ch(E_1 \otimes E_2) &= ch(E_1) \wedge ch(E_2). \end{aligned} \tag{5}$$

Sometime we encounter other characteristic classes, such as Todd classes, Hirzebruch  $L$ -polynomials, and  $\hat{A}$  polynomials. They correspond to different basis' of invariant polynomials. To describe these, use eigenvalues  $x_1, \dots, x_k$  of  $\frac{\sqrt{-1}F}{2\pi}$ . For example, the total Chern classes can be expressed as

$$c(F) = \det \left( 1 + \frac{\sqrt{-1}F}{2\pi} \right) = \prod_{i=1}^k (1 + x_i).$$

(Tangentially, it is interesting to note that the right-hand side takes the form  $\prod_i c(L_i)$ , where  $L_i$  is a line bundle with a curvature given by  $x_i$  and  $c(L_i) = 1 + x_i$ . Thus, as far as the Chern classes are concerned, the vector bundle  $E$  behaves as a direct sum of the line bundles  $L_1 \oplus L_2 \oplus \dots \oplus L_k$ . This phenomenon is called the splitting principle.)

Using this notation, the Todd classes are defined by,

$$td(E) = \prod_i \frac{x_i}{1 - e^{-x_i}},$$

the  $L$ -polynomials are defined by,

$$L(E) = \prod_i \frac{x_i}{\tanh x_i},$$

and the  $\hat{A}$ -polynomials are defined by,

$$\hat{A}(E) = \prod_i \frac{x_i/2}{\sinh(x_i/2)}.$$

## Chern number

Remarkably, the Chern classes and the Chern characters are integral. That means that if we integrate, say,  $c_i(E)$  over any  $2i$ -cycle in  $M$  with integer coefficients, we find an integer that is independent of the choice of the connection of  $E$ . If  $2k \geq n$ , we can integrate  $c_n(F)$  over the entire manifold  $M$  and get the Chern number. Let us compute Chern numbers in some examples.

(1) Consider the monopole bundle over  $S^2$ . It has the  $U(1)$  gauge field  $A$ . Let us denote the northern and southern hemispheres of  $S^2$  as  $H_{\pm}$ , and the gauge fields on them as  $A_{\pm}$ . For the monopole bundle with  $n$  monopole charge, the gauge field transforms as

$$A_- = A_+ + nd\phi$$

across the equator, where  $\phi$  is the longitude of  $S^2$ . We can then evaluate the Chern number as

$$\begin{aligned} C_1 &= \int_{S^2} c_1 \\ &= \frac{-1}{2\pi} \left( \int_{H_+} A_+ + \int_{H_-} A_- \right) \\ &= \frac{-1}{2\pi} \int_{S^1} (A_+ - A_-) = \frac{1}{2\pi} \int_0^{2\pi} nd\phi = n. \end{aligned} \quad (6)$$

Thus, the monopole number is the first Chern number in this case.

(2) Consider an  $SU(2)$  bundle over  $S^4$ . We can then consider the second Chern number,

$$C_2 = \int_{S^4} c_2 = \frac{1}{8\pi^2} \int_{S^4} \text{tr} F \wedge F.$$

We again split  $S^4$  into  $H_{\pm}$  such that  $H_+ \cap H_- = S^3$ . Over  $S^3$ , the gauge field transforms as

$$A_- = \Omega^{-1} A_+ \Omega + \Omega^{-1} d\Omega.$$

When we integrate  $\text{tr} F_{\pm} \wedge F_{\pm}$  over  $H_{\pm}$ , we note that the integrand can be written as

$$\text{tr} F \wedge F = d\text{tr} \left( AdA + \frac{2}{3} A^3 \right).$$

Note that this does not mean that  $\text{tr} F \wedge F$  is an exact form since the right-hand side, called the Chern-Simons form, is not necessarily globally defined over  $S^4$ . Thus,

$$\begin{aligned} C_2 &= \int_{S^4} c_2 \\ &= \frac{1}{8\pi^2} \int_{S^3} \text{tr} \left( A_+ dA_+ + \frac{2}{3} A_+^3 - A_- dA_- - \frac{2}{3} A_-^3 \right) \\ &= \frac{1}{24\pi^2} \int_{S^3} \text{tr} (\Omega^{-1} d\Omega)^3. \end{aligned} \quad (7)$$

The gauge transformation matrix  $\Omega$  is a map from  $S^3$  to  $SU(2)$ . Since the group  $SU(2)$  is diffeomorphic to  $S^3$  as a manifold, we can think of it as a map from  $S^3$  to itself. Such a map

can be classified by its winding number, which turns out to be the same as the second Chern number in the above.

### Euler class, Poltryagn classes

Let us turn to real vector bundles. Suppose  $X$  is a  $2r \times 2r$  real and anti-symmetrix matrix. In this case, in addition to  $\text{tr}$  and  $\det$ , we can consider one more way to construct an invariant polynomial. That is the Pfaffian,

$$Pf(X) = \frac{(-1)^r}{2^r r!} \epsilon^{i_1 j_1 i_2 j_2 \dots i_r j_r} X_{i_1 j_1} X_{i_2 j_2} \dots X_{i_r j_r}.$$

Note that, for antisymmetric matrices, the Pfaffian is a square root of the determinant,

$$\det X = Pf(X)^2.$$

If  $X$  is real and anti-symmetric, we can block-diagonalize it by  $SO(2r)$  as

$$X = \begin{pmatrix} 0 & x_1 & 0 & 0 & \dots & 0 & 0 \\ -x_1 & 0 & 0 & 0 & \dots & \cdot & 0 \\ 0 & 0 & 0 & x_2 & & & \\ 0 & 0 & -x_2 & 0 & & & \\ \cdot & \cdot & & & \cdot & & \\ 0 & \cdot & & & & 0 & x_r \\ 0 & 0 & & & & -x_r & 0 \end{pmatrix}$$

We can then write the Pfaffian as

$$Pf(X) = (-1)^r \prod_{i=1}^r x_i.$$

Under the conjugation  $X \rightarrow \Omega^t X \Omega$ , the Pfaffian transforms as

$$Pf(\Omega^t X \Omega) = \det \Omega \cdot Pf(X).$$

Thus, if  $\Omega \in SO(2r)$ , the Pfaffian is invariant. (Note that if  $\Omega \in O(2r)$ , the Pfaffian may change its sign.)

We can now define the Euler class by  $e(F) = Pf(F)$ .

In particular, the tangent bundle  $TM$  of an orientable Riemannian manifold  $M$  of dimensions  $n = 2r$  is an  $SO(2r)$  bundle. For example,

$$\begin{aligned} n = 2: \quad e(TM) &= \frac{1}{2\pi} R_{12}, \\ n = 4: \quad e(TM) &= \frac{1}{32\pi^2} \epsilon_{\mu\nu\rho\sigma} R^{\mu\nu} \wedge R^{\rho\sigma}, \end{aligned} \tag{8}$$

where the Riemann curvature is regarded as the 2-form as,

$$R^a_b = \frac{1}{2} R_{cd}^a e^c \wedge e^d = \frac{1}{2} R_{\mu\nu}^a dx^\mu \wedge dx^\nu.$$

The Gauss-Bonnet theorem for an even-dimensional manifold  $M$  relates the Euler characteristic  $\chi(M) = \sum_p (-1)^p b_p$  to the Euler class by

$$\chi(M) = \int_M e(TM).$$

The Pontryagin classes are defined similarly to the Chern classes as

$$p(E) = 1 + p_1(E) + p_2(E) + \cdots = \det \left( 1 - \frac{1}{2\pi} F \right).$$

Since  $F$  is anti-symmetric, we only have nontrivial polynomials with even degrees in  $F$ . Thus, we choose  $p_i(E)$  to be a  $4i$ -form. The highest Pontryagin class is at  $i = r$  where  $2r$  is the dimension of the fiber (unless  $2r > n = \dim M$ ). At this highest degree, it is the square of the Euler class,

$$p_r(E) = e(E)^2.$$