## Nonlinear Finite Element Method

04/10/2004

## Nonlinear Finite Element Method

- Lectures include discussion of the nonlinear finite element method.
- It is preferable to have completed "Introduction to Nonlinear Finite Element Analysis" available in summer session.
- If not, students are required to study on their own before participating this course. Reference:Toshiaki.,Kubo. "Introduction: Tensor Analysis For Nonlinear Finite Element Method" (Hisennkei Yugen Yoso no tameno Tensor Kaiseki no Kiso),Maruzen.
- Lecture references are available and downloadable at <a href="http://www.sml.k.utokyo.ac.jp/members/nabe/lecture2004">http://www.sml.k.utokyo.ac.jp/members/nabe/lecture2004</a> They should be put up on the website by the day before scheduled meeting day, and each students are expected to come in with a copy of the reference.
- •Lecture notes from previous year are available and downloadable, also at http://www.sml.k.u.tokyo.ac.jp/members/nabe/lecture2003 You may find the course title, "Advanced Finite Element Method" but the contents covered are almost the same I will cover this year.
- I will assign the exercises from this year, and expect the students to hand them in during the following lecture. They are not the requirements and they will not be graded, however it is important to actually practice calculate in deeper understanding the finite element method.
- For any questions, contact me at nabe@sml.k.u-tokyo.ac.jp

### Nonlinear Finite Element Method Lecture Schedule

- 1. 10/4 Finite element analysis for boundary value problem in differential equation
- 2. 10/18 Finite element analysis for linear elastic body
- 3. 10/25 Isoparametric solid element (program)
- 4. 11/1 Numerical solution and boundary condition processing for system of linear equations (with exercises)
- 5. 11/8 Basic program structure of the linear finite element method(program)
- 6. 11/15 Finite element formulation for geometric nonlinear problems(program)
- 7. 11/22 Static analysis technique, hyperelastic body and elastic-plastic material for nonlinear equations (program)
- 8. 11/29 Exercises for Lecture7
- 9. 12/6 Dynamic analysis technique and eigenvalue analysis of the nonlinear equations
- 10. 12/13 Structural element
- 11. 12/20 Numerical solution—skyline method, iterative method for the system of linear equations
- 12. 1/17 ALE finite element fluid analysis
- 13. 1/24 ALE finite element fluid analysis

## Review on the Finite Element Method — For Typical Numerical Solutions

- In physics, various phenomenon are interpreted in mathematical representations (differential equations, for example, are used in Newtonian mechanics).
- Finite element analysis is an approximate analysis for differential equations
- Finite difference calculus:
- Directly discretize the governing equations, and is theoretically simple, and it is easy to program.
- Which is met for Euler representation, however, not suitable for Lagrange representation.
- Finite element method:
- Discretize the governing equation in weak formulation
- Common analysis for solid-structural problems and fluid
- Programming is complicated

## Boundary Value Problems in Differential Equations

• Consider the boundary value problems in differential equations. Obtain u that satisfies the conditions for [B]

$$-\frac{d^2u}{dx^2} = f(x) \qquad (0 < x < a) \tag{1}$$

$$u(0) = u_0, \quad u(a) = u_a$$
 (2)

- •If f(x) is concerned with the problems that are analytically integrable, the solution u can be obtained by direct integration.
- Finite element method involves with approximate analysis of the weak form of the differential equations. The processes of obtaining the approximate solution are:
- Weak formulation of the differential equations.
- Eliminate the finite element interpolation function to substitute the continuous functions into the discrete nodal point value.
- Based on the above, construct the system of linear equations to obtain the approximate solution.

## Strong Form and Weak Form

• Weak form: obtain u that satisfies the conditions for [V].  $\int_0^a \frac{du \, dv}{dx \, dx} \, dx = \int_0^a f(x) \, v \, dx \quad \forall v \tag{3}$ 

$$u(0) = u_0, \quad u(a) = u_a \tag{4}$$

v is the function, which satisfies v(0) = v(a) = 0 also, the second order differentiable.

 Strong form: Obtain u that satisfies the conditions for [B].

$$-\frac{d^2u}{dx^2} = f(x) \qquad (0 < x < a) \tag{5}$$

$$u(0) = u_0, \quad u(a) = u_a$$
 (6)

# Deriving the Weak Formulation of the Differential Equations 1

Given the differential equation:

$$-\frac{d^2u}{dx^2} = f(x) \qquad (0 < x < a)$$
 (7)

Multiply the second-order differentiable function v, which satisfies the condition v(0) = v(a) = 0, on both sides then conduct integration of domain(0 to a).  $-\int_a^a \frac{d^2u}{dx^2} v \, dx = \int_a^a f(x) \, v \, dx \tag{8}$ 

• Integrate the left hand side equation in the above by parts.

$$\frac{d}{dx}\left(\frac{du}{dx}v\right) = \frac{d^2u}{dx^2}v + \frac{du}{dx}\frac{dv}{dx} \tag{9}$$

So,

$$-\frac{d^2u}{dx^2}v = -\frac{d}{dx}\left(\frac{du}{dx}v\right) + \frac{du}{dx}\frac{dv}{dx}$$
(10)

Given that,

$$\int_0^a \frac{du}{dx} \frac{dv}{dx} dx - \left[ \frac{du}{dx} v \right]_0^a = \int_0^a f(x) v dx \tag{11}$$

# Deriving the Weak Formulation of the Differential Equations 2

• Given, v(0) = v(a) = 0 the second term in the left hand side becomes 0

$$\int_0^a \frac{du}{dx} \frac{dv}{dx} dx - \left[ \frac{du}{dx} v \right]_0^a = \int_0^a f(x) v dx \tag{12}$$

Thus

$$\int_0^a \frac{du\,dv}{dx}\,dx = \int_0^a f(x)\,v\,dx\tag{13}$$

• By satisfying the boundary condition v(0) = v(a) = 0, we can also establish for the arbitrary v. Hence, [V] is proofed from [B].

# Deriving the Weak Formulation of the Differential Equations 3

Moreover, by tracing back the proof

$$\int_0^a \frac{du}{dx} \frac{dv}{dx} dx = \int_0^a f(x) v dx \tag{14}$$

v that satisfies the above must also satisfy the following equation.

$$-\int_{0}^{a} \frac{d^{2}u}{dx^{2}} v \, dx = \int_{0}^{a} f(x) \, v \, dx \tag{15}$$

Therefore,

$$\int_0^a \left\{ \frac{d^2 u}{dx^2} + f(x) \right\} v \, dx = 0 \tag{16}$$

In order to verify this for the arbitrary  $\nu$ , there is a necessary and sufficient condition:

$$\frac{d^2u}{dx^2} + f(x) = 0\tag{17}$$

• It is proofed that the solution *u* of [*V*] does not exist in any other but only in the solution of [*B*].

## Strong Form and Weak Form

• Therefore, [B] is equivalent with [V]. Find u that satisfies the conditions for [V].

$$\int_0^a \frac{du}{dx} \frac{dv}{dx} dx = \int_0^a f(x) v dx \quad \forall v$$
 (18)

Find u that satisfies the conditions for [B].

$$-\frac{d^2u}{dx^2} = f(x) \qquad (0 < x < a)$$

$$u(0) = u_0, \quad u(a) = u_a$$
(19)

• [V] is called the weak formulation of the differential equations.

• Furthermore, this [V] is equivalent with the following minimization problem [M]

Find 
$$u$$
 that satisfies the condition for  $UI$   $\int_0^a \frac{du}{dx} \frac{dv}{dx} dx = \int_0^a f(x) v dx \quad \forall v$  (21)

Find 
$$u$$
 that satisfies the condition for  $M$  (22)

Provided that, 
$$W(v) = \int_0^a \left(\frac{1}{2}\frac{dv}{dx}\frac{dv}{dx} - f(x)v\right)dx$$
 (23)

- In reality, in conducting finite analysis formulation of the solid structural analysis, such minimization problems are often used. (more accurately, stationary problems)
- This is because,
- There are some material models, which may be defined by such potentials.
- The matrix obtained is always found in symmetry.
- Lagrange method of undetermined multipliers can be applied.

• To show  $[M] \Rightarrow [V]$ , if we define the equation above to be v = u + w then, which becomes equivalent with the following equation:

$$W(u+w) - W(u) \ge 0 \qquad \forall w \tag{24}$$

• Here, write out the equation based on the definition for W(v):

$$\int_{0}^{a} \left[ \left\{ \frac{1}{2} \left( \frac{d(u+w)}{dx} \frac{d(u+w)}{dx} \right) - f(x) (u+w) \right\} - \left\{ \frac{1}{2} \frac{du}{dx} \frac{du}{dx} - f(x) u \right\} \right] dx$$

$$= \int_{0}^{a} \left[ \left\{ \frac{du}{dx} \frac{dw}{dx} - f(x) w \right\} + \frac{1}{2} \frac{dw}{dx} \frac{dw}{dx} \right] dx \ge 0$$
(25)

• This equation must be also established for the arbitrary w, thus we define  $\alpha$  as the arbitrary scalar to substitute  $\alpha w$  into w. To put it in order:

$$\alpha \int_0^a \left\{ \frac{du}{dx} \frac{dw}{dx} - f(x) w \right\} dx + \frac{\alpha^2}{2} \int_0^a \frac{dw}{dx} \frac{dw}{dx} dx \ge 0$$
 (26)

- •Where  $w \equiv 0$  ,the left hand side becomes 0 thus it is established.
- •Where  $w \not\equiv 0$  ,we consider it as the quadratic equations of  $\alpha$ .
- Therefore,

$$y = a\alpha^{2} + b\alpha + c$$

$$a = \frac{1}{2} \int_{0}^{a} \frac{dw}{dx} \frac{dw}{dx} dx$$

$$b = \int_{0}^{a} \left\{ \frac{du}{dx} \frac{dw}{dx} - f(x) w \right\} dx$$

$$c = 0$$
(27)

• A necessary and sufficient condition for  $y \ge 0$  is where a > 0 and  $b^2 - 4ac \le 0$  (discriminant), however in this case the condition may be b = 0

•In other words, 
$$\int_0^a \left\{ \frac{du \, dw}{dx \, dx} - f(x) \, w \right\} dx = 0 \tag{28}$$

- •Thus, [V] should be proofed by [M].
- For  $[V] \Rightarrow [M]$ , we can make a proof by having a definition W(u+w)-W(u) to calculate then gain the following:

$$W(u+w) - W(u)$$

$$= \int_0^a \left[ \left\{ \frac{du \, dw}{dx \, dx} - f(x) \, w \right\} + \frac{1}{2} \frac{dw \, dw}{dx} \right] dx$$

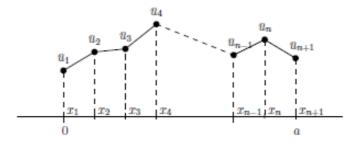
$$= \int_0^a \frac{1}{2} \frac{dw \, dw}{dx} dx \ge 0 \tag{29}$$

### Approximate Solutions for Weak Form— Dividing the Interval of Integrations

- Obtain the approximate solution for the weak form by the finite element method.
- First, divide the domain [0, a] of u into the n intervals of [xi, xi+1]  $(i = 1 \sim n)$  where there is no overlapping of u.
- This xi ( $i = 1 \sim n + 1$ ) is called a nodal point.
- Now the integration of  $\int_0^a \frac{du}{dx} \frac{dv}{dx} dx = \int_0^a f(x) v dx$  can be reformed as the following.

$$\sum_{i=1}^{n} \int_{x_i}^{x_{i+1}} \frac{du}{dx} \frac{dv}{dx} dx = \sum_{i=1}^{n} \int_{x_i}^{x_{i+1}} f \cdot v \, dx \tag{30}$$

• We assume the approximate solution to take the value ui ( $i = 1 \sim n + 1$ ) in the nodal point xi ( $i = 1 \sim n + 1$ ), and where in between the nodal points xi and xi+1, we assume the solution takes changes linearly. In this case, we use the similar functions of v.



## Finite Element Method Interpolation1

• In a convenient means, we express each interval of integration, the initial point xi and the end point xi+1 as x(i,1) and x(i,2) respectively.

$$\sum_{i=1}^{n} \int_{x_{(i,1)}}^{x_{(i,2)}} \frac{du}{dx} \frac{dv}{dx} dx = \sum_{i=1}^{n} \int_{x_{(i,1)}}^{x_{(i,2)}} f \cdot v \, dx \tag{31}$$

• Conduct change of variables in order for x(i,1) to correspond to -1, and for x(i,2) to correspond to 1 in each intervals.

$$u_{1}$$
 $u_{2}$ 
 $u_{3}$ 
 $u_{n-1}$ 
 $u_{n+1}$ 
 $u_{n+1}$ 

## Finite Element Method Interpolation2

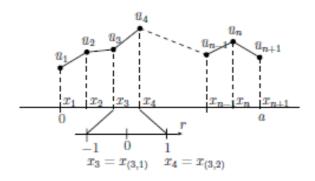
• Now, using the parameter r ( $-1 \le r \le 1$ ), x found in [x(i,1), x(i,2)] can be expressed in the following form.

$$x = N^{(1)}x_{(i,1)} + N^{(2)}x_{(i,2)}$$
 (32)

$$N^{(1)} = \frac{1}{2}(1-r) \tag{33}$$

$$N^{(2)} = \frac{1}{2}(1+r) \tag{34}$$

• N(i) is called the finite element method interpolation function, or simply the interpolation function.



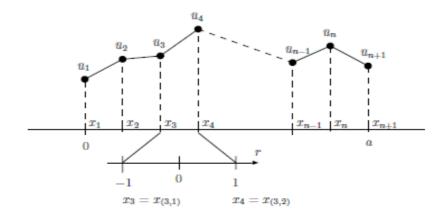
## Finite Element Method Interpolation3

• Likewise, for  $\bar{u}, \bar{v}$ , using the parameter  $r(-1 \le r \le 1)$  to express in the following form.

$$\bar{u} = N^{(1)}\bar{u}_{(i,1)} + N^{(2)}\bar{u}_{(i,2)}$$
 (35)

$$\bar{v} = N^{(1)}\bar{v}_{(i,1)} + N^{(2)}\bar{v}_{(i,2)}$$
 (36)

Only where,  $\bar{u}_{(i,1)} = \bar{u}_i$ ,  $\bar{u}_{(i,2)} = \bar{u}_{i+1}$ ,  $\bar{v}_{(i,1)} = \bar{v}_i$ ,  $\bar{v}_{(i,2)} = \bar{v}_{i+1}$ 



## Discrete Value Expression of Differentiation

Substitute the weak form into the approximate solution to obtain the following:

$$\sum_{i=1}^{n} \int_{x_{(i,1)}}^{x_{(i,2)}} \frac{d\bar{u}}{dx} \frac{d\bar{v}}{dx} dx = \sum_{i=1}^{n} \int_{x_{(i,1)}}^{x_{(i,2)}} f \cdot \bar{v} dx$$
(37)

In the left hand side, the derivative of the function x of  $\bar{u}$ ,  $\bar{v}$  in the integrand can be obtained by using the chain rule.

$$\frac{d\bar{u}}{dx} = \frac{d}{dx} \left( N^{(1)} \bar{u}_{(i,1)} + N^{(2)} \bar{u}_{(i,2)} \right) 
= \frac{dN^{(1)}}{dx} \bar{u}_{(i,1)} + \frac{dN^{(2)}}{dx} \bar{u}_{(i,2)} 
= \frac{dN^{(1)}}{dr} \frac{dr}{dx} \bar{u}_{(i,1)} + \frac{dN^{(2)}}{dr} \frac{dr}{dx} \bar{u}_{(i,2)}$$
(38)

•For  $\overline{dx}$  appears in the equation above, the inverse numbers may be obtained. In actual programming,

We obtain  $\frac{dr}{dx}$  in the beginning then substitute it into the above equation to follow the calculations shows in below:

$$\frac{dx}{dr} = \frac{d}{dr} \left( N^{(1)} x_{(i,1)} + N^{(2)} x_{(i,2)} \right) 
= \frac{dN^{(1)}}{dr} x_{(i,1)} + \frac{dN^{(2)}}{dr} x_{(i,2)}$$
(39)

### **Element Matrix**1

• The Approximation of the weak form:

$$\sum_{i=1}^{n} \int_{x_{(i,1)}}^{x_{(i,2)}} \frac{d\bar{u}}{dx} \frac{d\bar{v}}{dx} dx = \sum_{i=1}^{n} \int_{x_{(i,1)}}^{x_{(i,2)}} f \cdot \bar{v} dx$$
(40)

The left hand side integrand can be expressed in the following matrix:

$$\frac{d\bar{u}}{dx}\frac{d\bar{v}}{dx} = \left(\frac{dN^{(1)}}{dx}\bar{u}_{(i,1)} + \frac{dN^{(2)}}{dx}\bar{u}_{(i,2)}\right) \left(\frac{dN^{(1)}}{dx}\bar{v}_{(i,1)} + \frac{dN^{(2)}}{dx}\bar{v}_{(i,2)}\right) \\
= \left\{\bar{v}_{(i,1)}\bar{v}_{(i,2)}\right\} \begin{bmatrix} \frac{dN^{(1)}}{dx} \\ \frac{dN^{(2)}}{dx} \end{bmatrix} \begin{bmatrix} \frac{dN^{(1)}}{dx}\frac{dN^{(2)}}{dx} \end{bmatrix} \begin{cases} \bar{u}_{(i,1)} \\ \bar{u}_{(i,2)} \end{cases} \\
= \left\{\bar{v}_{(i,1)}\bar{v}_{(i,2)}\right\} \begin{bmatrix} \frac{dN^{(1)}}{dx}\frac{dN^{(1)}}{dx}\frac{dN^{(1)}}{dx}\frac{dN^{(2)}}{dx}\frac{dN^{(2)}}{dx} \end{bmatrix} \begin{cases} \bar{u}_{(i,1)} \\ \bar{u}_{(i,2)} \end{cases}$$

$$(41)$$

• Suppose, we have  $J_{(i)} = \frac{dx}{dr}$  in each interval of integration [x(i,1), x(i,2)] then, we can write as following:

$$\sum_{i=1}^{n} \int_{-1}^{1} \{\bar{v}_{(i,1)} \, \bar{v}_{(i,2)}\} \begin{bmatrix} \frac{dN^{(1)}}{dx} \frac{dN^{(1)}}{dx} & \frac{dN^{(1)}}{dx} \frac{dN^{(2)}}{dx} \\ \frac{dN^{(2)}}{dx} \frac{dN^{(1)}}{dx} & \frac{dN^{(2)}}{dx} \frac{dN^{(2)}}{dx} \end{bmatrix} \begin{bmatrix} \bar{u}_{(i,1)} \\ \bar{u}_{(i,2)} \end{bmatrix} J_{(i)} \, dr = \sum_{i=1}^{n} \int_{-1}^{1} \{\bar{v}_{(i,1)} \, \bar{v}_{(i,2)}\} \begin{bmatrix} N^{(1)} \\ N^{(2)} \end{bmatrix} f \, J_{(i)} \, dr$$
(42)

• The following equation may be obtained with this tedious work of writing out the matrix. It is important to realize dx/dr and dr/dx are used in a different meaning.

$$\sum_{i=1}^{n} \int_{-1}^{1} \{\bar{v}_{(i,1)} \, \bar{v}_{(i,2)}\} \begin{bmatrix} \frac{dN^{(1)}}{dr} \frac{dr}{dx} \frac{dN^{(1)}}{dr} \frac{dr}{dx} & \frac{dN^{(1)}}{dr} \frac{dr}{dx} \frac{dN^{(2)}}{dr} \frac{dr}{dx} \\ \frac{dN^{(2)}}{dr} \frac{dr}{dx} \frac{dN^{(1)}}{dr} \frac{dr}{dx} & \frac{dN^{(2)}}{dr} \frac{dr}{dx} \frac{dN^{(2)}}{dr} \frac{dr}{dx} \end{bmatrix} \begin{bmatrix} \bar{u}_{(i,1)} \\ \bar{u}_{(i,2)} \end{bmatrix} \frac{dx}{dr} dr = \sum_{i=1}^{n} \int_{-1}^{1} \{\bar{v}_{(i,1)} \, \bar{v}_{(i,2)}\} \begin{cases} N^{(1)} \\ N^{(2)} \end{cases} f \frac{dx}{dr} dr$$

$$(43)$$

### Element Matrix 2

•  $\bar{v}_{(i,1)}$ ,  $\bar{v}_{(i,2)}$ ,  $\bar{u}_{(i,1)}$ ,  $\bar{u}_{(i,2)}$  are the values in the nodal points, and remain constant in terms of r, the variable of integration, therefore we may leave them out of the integrals.

$$\sum_{i=1}^{n} \int_{-1}^{1} \{\bar{v}_{(i,1)} \, \bar{v}_{(i,2)}\} \left\{ \frac{dN^{(1)}}{dx} \frac{dN^{(1)}}{dx} \frac{dN^{(1)}}{dx} \frac{dN^{(2)}}{dx} \frac{dN^{(2)}}{dx} \right\} \left\{ \bar{u}_{(i,1)} \right\} J_{(i)} dr$$

$$= \sum_{i=1}^{n} \{\bar{v}_{(i,1)} \, \bar{v}_{(i,2)}\} \int_{-1}^{1} \left[ \frac{dN^{(1)}}{dx} \frac{dN^{(1)}}{dx} \frac{dN^{(1)}}{dx} \frac{dN^{(2)}}{dx} \frac{dN^{(2)}}{dx} \right] J_{(i)} dr \left\{ \bar{u}_{(i,1)} \right\} \left\{ \bar{u}_{(i,2)} \right\} \left\{ \frac{dN^{(2)}}{dx} \frac{dN^{(1)}}{dx} \frac{dN^{(2)}}{dx} \frac{dN^{(2)}}{dx} \frac{dN^{(2)}}{dx} \right\} J_{(i)} dr \left\{ \bar{u}_{(i,1)} \right\} \left\{ \bar{u}_{(i,2)} \right\} \left\{ N^{(1)} \right\} f J_{(i)} dr = \sum_{i=1}^{n} \{\bar{v}_{(i,1)} \, \bar{v}_{(i,2)}\} \int_{-1}^{1} \left\{ N^{(1)} \right\} f J_{(i)} dr$$

$$(45)$$

• [K(i)], which defines  $\{F(i)\}$  in the following equation, is called the element matrix.

$$[K^{(i)}] = \begin{bmatrix} K_{11}^{(i)} & K_{12}^{(i)} \\ K_{21}^{(i)} & K_{22}^{(i)} \end{bmatrix} = \int_{-1}^{1} \begin{bmatrix} \frac{dN^{(1)}}{dx} \frac{dN^{(1)}}{dx} & \frac{dN^{(1)}}{dx} \frac{dN^{(2)}}{dx} \\ \frac{dN^{(1)}}{dx} \frac{dN^{(2)}}{dx} & \frac{dN^{(2)}}{dx} \frac{dN^{(2)}}{dx} \end{bmatrix} J_{(i)} dr \qquad \{F^{(i)}\} = \begin{cases} F_{1}^{(i)} \\ F_{2}^{(i)} \end{cases} = \int_{-1}^{1} \begin{cases} N^{(1)} \\ N^{(2)} \end{cases} f J_{(i)} dr$$

$$(46)$$

ullet Write out the both sides of the equation without using  $\Sigma$ .

$$\sum_{i=1}^{n} \{\bar{v}_{(i,1)} \, \bar{v}_{(i,2)}\} \begin{vmatrix} K_{11}^{(i)} & K_{12}^{(i)} \\ K_{21}^{(i)} & K_{22}^{(i)} \end{vmatrix} \begin{cases} \bar{u}_{(i,1)} \\ \bar{u}_{(i,2)} \end{pmatrix} = \sum_{i=1}^{n} \{\bar{v}_{(i,1)} \, \bar{v}_{(i,2)}\} \begin{cases} F_{1}^{(i)} \\ F_{2}^{(i)} \end{cases}$$

$$(47)$$

Therefore,

$$\sum_{i=1}^{n} \{\bar{v}_{(i,1)} \, \bar{v}_{(i,2)}\} \begin{bmatrix} K_{11}^{(i)} & K_{12}^{(i)} \\ K_{21}^{(i)} & K_{22}^{(i)} \end{bmatrix} \begin{cases} \bar{u}_{(i,1)} \\ \bar{u}_{(i,2)} \end{cases} 
= \{\bar{v}_{(1,1)} \, \bar{v}_{(1,2)}\} \begin{bmatrix} K_{11}^{(1)} & K_{12}^{(1)} \\ K_{21}^{(1)} & K_{22}^{(1)} \end{bmatrix} \begin{cases} \bar{u}_{(1,1)} \\ \bar{u}_{(1,2)} \end{pmatrix} + \{\bar{v}_{(2,1)} \, \bar{v}_{(2,2)}\} \begin{bmatrix} K_{11}^{(2)} & K_{12}^{(2)} \\ K_{21}^{(2)} & K_{22}^{(2)} \end{bmatrix} \begin{cases} \bar{u}_{(2,1)} \\ \bar{u}_{(2,2)} \end{pmatrix} 
+ \dots + \{\bar{v}_{(n,1)} \, \bar{v}_{(n,2)}\} \begin{bmatrix} K_{11}^{(n)} & K_{12}^{(n)} \\ K_{21}^{(n)} & K_{22}^{(n)} \end{bmatrix} \begin{cases} \bar{u}_{(n,1)} \\ \bar{u}_{(n,2)} \end{pmatrix}$$
(48)

$$\sum_{i=1}^{n} \{\bar{v}_{(i,1)} \, \bar{v}_{(i,2)}\} \begin{Bmatrix} F_1^{(i)} \\ F_2^{(i)} \end{Bmatrix} 
= \{\bar{v}_{(1,1)} \, \bar{v}_{(1,2)}\} \begin{Bmatrix} F_1^{(1)} \\ F_2^{(1)} \end{Bmatrix} + \{\bar{v}_{(2,1)} \, \bar{v}_{(2,2)}\} \begin{Bmatrix} F_1^{(2)} \\ F_2^{(2)} \end{Bmatrix} + \dots + \{\bar{v}_{(n,1)} \, \bar{v}_{(n,2)}\} \begin{Bmatrix} F_1^{(n)} \\ F_2^{(n)} \end{Bmatrix}$$
(49)

• Whereas,  $\bar{u}_1 = \bar{u}_{(i,1)}$ ,  $\bar{u}_2 = \bar{u}_{(i,2)}$ ,  $\bar{v}_1 = \bar{v}_{(i,1)}$ ,  $\bar{v}_2 = \bar{v}_{(i,2)}$  we can force to write in the matrix representation shows in the following:

$$\{\bar{v}_{(1,1)}\,\bar{v}_{(1,2)}\}\begin{bmatrix}K_{11}^{(1)} & K_{12}^{(1)} \\ K_{21}^{(1)} & K_{22}^{(1)}\end{bmatrix}\begin{bmatrix}\bar{u}_{(1,1)} \\ \bar{u}_{(1,2)}\end{bmatrix}\} = \{\bar{v}_1\,\bar{v}_2\}\begin{bmatrix}K_{11}^{(1)} & K_{12}^{(1)} \\ K_{21}^{(1)} & K_{22}^{(1)}\end{bmatrix}\begin{bmatrix}\bar{u}_1 \\ \bar{u}_2\end{bmatrix} = \{\bar{v}_1,\bar{v}_2,\ldots,\bar{v}_{n+1}\}\begin{bmatrix}K_{11}^{(1)} & K_{12}^{(1)} & 0 & \ldots & 0 \\ K_{21}^{(1)} & K_{22}^{(1)} & 0 & \ldots & 0 \\ 0 & 0 & 0 & \ldots & 0 \\ \vdots & \vdots & \vdots & \vdots & \vdots \\ 0 & 0 & 0 & 0\end{bmatrix}\begin{bmatrix}\bar{u}_1 \\ \bar{u}_2 \\ \vdots \\ \bar{u}_{n+1}\end{bmatrix} \tag{50}$$

The second matrix may be written as:

$$\begin{aligned}
&\{\bar{v}_{(2,1)}\,\bar{v}_{(2,2)}\} \begin{bmatrix} K_{11}^{(2)} & K_{12}^{(2)} \\ K_{21}^{(2)} & K_{22}^{(2)} \end{bmatrix} \begin{cases} \bar{u}_{(2,1)} \\ \bar{u}_{(2,2)} \end{pmatrix} = \{\bar{v}_2\,\bar{v}_3\} \begin{bmatrix} K_{11}^{(2)} & K_{12}^{(2)} \\ K_{21}^{(2)} & K_{22}^{(2)} \end{bmatrix} \begin{cases} \bar{u}_2 \\ \bar{u}_3 \end{pmatrix} \\
&= \{\bar{v}_1, \bar{v}_2, \dots, \bar{v}_{n+1}\} \begin{bmatrix} 0 & 0 & 0 & 0 & \cdots & 0 \\ 0 & K_{11}^{(2)} & K_{12}^{(2)} & 0 & \cdots & 0 \\ 0 & K_{21}^{(2)} & K_{22}^{(2)} & 0 & \cdots & 0 \\ 0 & K_{21}^{(2)} & K_{22}^{(2)} & 0 & \cdots & 0 \\ 0 & \vdots & \vdots & \vdots & \vdots \\ 0 & 0 & 0 & 0 & 0 \end{bmatrix} \begin{cases} \bar{u}_1 \\ \bar{u}_2 \\ \vdots \\ \bar{u}_{n+1} \end{aligned} \right\} \tag{51}$$

• In the same way, for the *i*th matrix, we can generalize to write as the following:

$$\{\bar{v}_{(i,1)}\,\bar{v}_{(i,2)}\} \begin{bmatrix} K_{11}^{(i)} & K_{12}^{(i)} \\ K_{21}^{(i)} & K_{22}^{(i)} \end{bmatrix} \begin{Bmatrix} \bar{u}_{(i,1)} \\ \bar{u}_{(i,2)} \end{Bmatrix} = \{\bar{v}_1, \bar{v}_2, \dots, \bar{v}_{n+1}\} \begin{bmatrix} 0 & \cdots & 0 & 0 & 0 & \cdots & 0 \\ 0 & \cdots & \vdots & \vdots & \vdots & \vdots \\ 0 & \cdots & 0 & K_{11}^{(i)} & K_{12}^{(i)} & 0 & \cdots & 0 \\ 0 & \cdots & 0 & K_{21}^{(i)} & K_{22}^{(i)} & 0 & \cdots & 0 \\ 0 & \cdots & \vdots & \vdots & \vdots & \vdots \\ 0 & \cdots & 0 & 0 & 0 & 0 & 0 \end{bmatrix} \begin{Bmatrix} \bar{u}_1 \\ \bar{u}_2 \\ \vdots \\ \bar{u}_{n+1} \end{Bmatrix}$$
(52)

Based on the above,

$$\sum_{i=1}^{n} \{\bar{v}_{(i,1)} \, \bar{v}_{(i,2)}\} \begin{bmatrix} K_{11}^{(i)} & K_{12}^{(i)} \\ K_{21}^{(i)} & K_{22}^{(i)} \end{bmatrix} \begin{cases} \bar{u}_{(i,1)} \\ \bar{u}_{(i,2)} \end{pmatrix} \\
= \{\bar{v}_{1}, \bar{v}_{2}, \dots, \bar{v}_{n+1}\} \begin{bmatrix} K_{11}^{(1)} & K_{12}^{(1)} \\ K_{21}^{(1)} & K_{22}^{(1)} + K_{11}^{(2)} & K_{12}^{(2)} \\ K_{21}^{(2)} & K_{22}^{(2)} + K_{11}^{(3)} & K_{12}^{(2)} \\ & & & & & \\ & & & & & \\ & & & & & \\ & & & & & \\ & & & & & \\ \end{bmatrix} \begin{bmatrix} \bar{u}_{1} \\ \bar{u}_{2} \\ \vdots \\ \bar{u}_{n+1} \end{bmatrix} \tag{53}$$

• The  $n+1 \times n+1$  dimensional matrix [K] made by superposition of the element matrix is called the compound matrix.

$$[K] = \begin{bmatrix} K_{11}^{(1)} & K_{12}^{(1)} \\ K_{21}^{(1)} & K_{22}^{(1)} + K_{11}^{(2)} & K_{12}^{(2)} \\ & K_{21}^{(2)} & K_{22}^{(2)} + K_{11}^{(3)} \\ & & K_{21}^{(3)} & & & & \\ & & & & & \\ & & & & & & \\ & & & & & \\ & & & & & \\ & & & & & \\ & & & & & \\ & & & & & \\ & & & & & \\ & & & & & \\ & & & & & \\ & & & & & \\ & & & & & \\ & & & & & \\ & & & & & \\ & & & & & \\ & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & \\ & & & & \\ & & & \\ & & & & \\ &$$

• The component of  $\{F\}$  can be written as.

$$\{F\} = \begin{cases} F_1^{(1)} \\ F_2^{(1)} + F_1^{(2)} \\ F_2^{(2)} + F_1^{(3)} \\ \vdots \end{cases}$$
 (55)

• Using this matrix notation in the above, the following equation (56)can be written as (57)

$$\sum_{i=1}^{n} \int_{x_{(i,1)}}^{x_{(i,2)}} \frac{d\bar{u}}{dx} \frac{d\bar{v}}{dx} dx = \sum_{i=1}^{n} \int_{x_{(i,1)}}^{x_{(i,2)}} f \cdot \bar{v} dx \qquad (56)$$

$$\{\bar{v}\}^T[K]\{\bar{u}\} = \{\bar{v}\}^T\{F\}$$
 (57)

• The following condition is required for the arbitrary v

$$\{\bar{v}\}^T[K]\{\bar{u}\} = \{\bar{v}\}^T\{F\}$$
 (58)

$$\{\bar{v}\}^T[[K]\{\bar{u}\} - \{F\}] = 0 \tag{59}$$

However, the condition above is established only when,

$$[K]\{\bar{u}\} = \{F\}$$
 (60)

• Hence the approximate solution is required for the weak form of the differential equations as the system of linear equations shows in the following,  $[K]\{\bar{u}\}=\{F\}$  (61)

• These are the basic procedures of the finite element method.

## Discretization Error of the Finite Element Method

- Finite element method is the approximate solution of the differential equations.
- While the solutions of the differential equations are continuous, the solutions of the finite element method are found within the range of unknown numbers.
- Essentially, involves the errors
- Which is called the discretization error
- Does error decrease in accordance with the finite element partitions becoming smaller?

## Terminology

Inner product: Assume V as linear space on R. When the bilinear form ( •, •):
 V × V → R satisfies the following condition, then it is called the inner product.

$$(u, v) = (v, u) \tag{62}$$

$$(cu, v) = c(u, v) \quad c \in \mathbb{R}$$
 (63)

$$(u + v, w) = (u, w) + (v, w)$$
 (64)

$$(u, u) \ge 0$$
  $(u, u) = 0 \Leftrightarrow u = 0$  (65)

• Norm : When  $\|\cdot\|:V\to R$  satisfies the following condition, then it is called norm.  $\|u\|\geq 0 \quad \|u\|=0 \Leftrightarrow u=0$  (66)

$$||cu|| = |c| ||u|| \quad c \in \mathbb{R}$$
 (67)

$$||u+v|| \le ||u|| + ||v||$$
 (68)

- • $(u, u)^{1/2}$ , for example, satisfy the condition for norm.
- Schwarz inequality

$$|(u, v)| \le ||u|| ||v||$$
 (69)

### Principle of Minimum Error 1

#### Organizing the symbols

$$(u,v)_H = \int_0^a \frac{du}{dx} \frac{dv}{dx} dx \tag{70}$$

$$(u,v)_L = \int_0^a u \, v \, dx \tag{71}$$

$$||u||_{H} = \{(u, u)_{H}\}^{(1/2)}$$
 (72)

$$||u||_{L} = \{(u, u)_{L}\}^{(1/2)} \tag{73}$$

#### Weak form

$$\int_0^a \frac{du}{dx} \frac{dv}{dx} dx = \int_0^a f(x) v dx \quad \Rightarrow \quad (u, v)_H = (f, v)_L \tag{74}$$

#### Discretization of the weak form

$$\int_0^a \frac{d\bar{u}}{dx} \frac{d\bar{v}}{dx} dx = \int_0^a f \,\bar{v} \,dx \quad \Rightarrow \quad (\bar{u}, \bar{v})_H = (f, \bar{v})_L \tag{75}$$

### Principle of Minimum Error 2

• We can take  $\bar{v}$  as v in (\*1)

$$(u,v)_H = (f,v)_L \tag{*1}$$

$$(\bar{u}, \bar{v})_H = (f, \bar{v})_L \tag{*2}$$

Thus

$$(u,\bar{v})_H = (f,\bar{v})_L \tag{*3}$$

•From (\*2) and (\*3),

$$(u - \bar{u}, \bar{v})_H = 0 \tag{*4}$$

• For the discretized solution candidate  $\bar{w}$ , we can take  $\bar{u} - \bar{w}$  as  $\bar{v}$  in (\*2).( $\bar{u} - \bar{w}$  becomes 0 at both extremes) According to(\*4),

$$(\bar{u} - u, \bar{u} - \bar{w})_H = 0 \tag{*5}$$

$$(\bar{u} - u, \bar{u} - \bar{w})_H = (\bar{u} - u, \bar{u} - u + u - \bar{w})_H = 0 \tag{*6}$$

$$(\bar{u} - u, \bar{u} - u)_H = (\bar{u} - u, \bar{w} - u)_H$$
 (\*7)

### Principle of Minimum Error 3

Using Schwarz inequality to the right hand side,

$$\|\bar{u} - u\|_{H}^{2} \le \|\bar{u} - u\|_{H} \|\bar{w} - u\|_{H}$$
 (76)

• So,

$$\| \bar{u} - u \|_{H} \le \| \bar{w} - u \|_{H}$$
 (77)

- The norm for  $\bar{u} u$  are consisted of an approximate solution, a correct solution, and some kind of error.
- ullet ar w may take the arbitrary solutions of all candidates.
- ullet ar u is the closest correct solution among all approximate solution candidates.
- This is called the principle of minimum error.

## Order of Magnitude Estimate and The Estimate of The Function Itself

 By taking the maximum length of the element h, we can make a proof for the following relation, and this is called the order of the magnitude estimate.

$$\| \bar{u} - u \|_{H} \le h \| f \|_{L}$$
 (78)

• For the error of the function itself  $\parallel \bar{u} - u \parallel_L$  we can proof the following relation:

$$\|\bar{u} - u\|_{L} \le h^{2} \|f\|_{L}$$
 (79)

 For further discussion, refer :Kikuchi.,Fumio."Yugen Yosoho Gaisetsu"

## Weak Form With More Common Boundary Condition

Boundary value problems of the differential equations

$$-\frac{d^2u}{dx^2} = f \qquad (0 < x < a) \tag{80}$$

$$u(0) = \alpha \quad \frac{du}{dx}(a) = \beta \tag{81}$$

• Suppose  $\check{u}$  satisfies the following. (Not necessarily(0) =  $\alpha$ )

$$-\frac{d^2\check{u}}{dx^2} = f \qquad (0 < x < a) \tag{82}$$

$$\frac{d\check{u}}{dx}(a) = \beta \tag{83}$$

• In respect, multiply v that satisfies  $v(0) = \alpha$ , then integrate.

$$-\int_0^a \frac{d^2\check{u}}{dx^2} v dx = \int_0^a f v dx \tag{84}$$

Here, using the integration by parts to reform the equation above.

$$\frac{d}{dx}\left(\frac{d\check{u}}{dx}v\right) = \frac{d^2\check{u}}{dx^2}v + \frac{d\check{u}}{dx}\frac{dv}{dx} \tag{85}$$

$$\frac{d^2\check{u}}{dx^2}v = \frac{d}{dx}\left(\frac{d\check{u}}{dx}v\right) - \frac{d\check{u}}{dx}\frac{dv}{dx} \tag{86}$$

$$-\int_0^a \frac{d^2 \check{u}}{dx^2} v dx = -\int_0^a \frac{d}{dx} \left(\frac{d\check{u}}{dx} v\right) dx + \int_0^a \frac{d\check{u}}{dx} \frac{dv}{dx} dx \tag{87}$$

$$= -\left[\frac{d\check{u}}{dx}v\right]_0^a + \int_0^a \frac{d\check{u}}{dx}\frac{dv}{dx}dx \tag{88}$$

$$= -\frac{d\check{u}}{dx}(a)v(a) + \frac{d\check{u}}{dx}(0)v(0) + \int_0^a \frac{d\check{u}}{dx}\frac{dv}{dx}dx$$
 (89)

(90)

$$-\frac{d\check{u}}{dx}(a)v(a) + \frac{d\check{u}}{dx}(0)v(0) + \int_0^a \frac{d\check{u}}{dx}\frac{dv}{dx}dx = \int_0^a fvdx \tag{91}$$

• Substitute the solution u as  $\check{u} \cdot v$  is decomposed in,

$$v = u + \delta u \tag{92}$$

Here  $v(0) = u(0) = \alpha$  is given,

$$\delta u(0) = 0 \tag{93}$$

$$-\frac{du}{dx}(a)\{u(a) + \delta u(a)\} + \frac{du}{dx}(0)\{u(0) + \delta u(0)\} + \int_0^a \frac{du}{dx} \left\{ \frac{du}{dx} + \frac{d\delta u}{dx} \right\} = \int_0^a f(u + \delta u) dx \quad (94)$$

In addition, if u is substituted for v,

$$-\frac{du}{dx}(a)u(a) + \frac{du}{dx}(0)u(0) + \int_0^a \frac{du}{dx}\frac{du}{dx}dx = \int_0^a fudx \tag{95}$$

Take the difference from each side,

$$-\underbrace{\frac{du}{dx}(a)}_{=\beta}\delta u(a) + \underbrace{\frac{du}{dx}(0)}_{=0}\underbrace{\delta u(0)}_{=0} + \int_0^a \frac{du}{dx}\frac{d\delta u}{dx}dx = \int_0^a f\delta udx \tag{96}$$

Then we obtain the following weak form,

$$\int_0^a \frac{du}{dx} \frac{d\delta u}{dx} dx = \int_0^a f \delta u dx + \beta \delta u(a) \tag{97}$$

## Method of Weighted Residuals

- In a further discussion, we cover the method of weighted residuals
- Given the governing equation in the following.

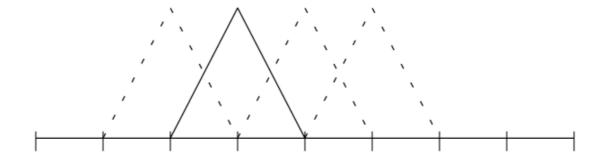
$$Q(u) = P(u) \tag{98}$$

- Suppose the approximate solution to be  $\bar{u} = Mr = 1$   $Cr \phi r$ , then the residuals  $R = \bar{u} Q(\bar{u} P(\bar{u}) P(\bar{u}) = 0$  can be obtained as  $\bar{R} = 0$   $\bar{u}$  hen  $\bar{u} = u$ . Now, obtain the function f(R) with its minimum value when wf(R) (f(R) = 0) when  $\bar{R} = 0$
- Define the linear independent function  $\psi q$  with orthogonal polynomials M, then orthogonalize the function with the residuals R.

$$\int_{v} \psi_q R dv = 0, \quad q = 1, \dots, M \tag{99}$$

•Galerkin method is a method of using  $\phi$  instead of  $\psi$ .

$$\int_{v} \phi_r R dv = 0, \quad r = 1, \dots, M \tag{100}$$



• As previous example shows,  $\phi_r$  forms the polygonal line taking the value1 at the nodal point r, and 0 in other nodal points.

## 2004 Advanced Nonlinear Finite Element Method: Exercises1

- Verify [V] from [B].(check the equations in the notes)
- Eq.(24) sets v = u + w, explain the reason why we can have such definition.
- Using Eq.(47), express the Eq.(47) (check the equations in the notes)
- With the region 0 < x < 1, and the number of partition n = 4, obtain the element matrix as well as the compound matrix. In obtaining the element matrix, remember to write out every matrix. Do not only write the last matrix obtained.
- Using Eq.(77), express the Eq.(97) (check the equation in the notes)
- On the nodal point r,  $\phi r$  takes the value 1, but for the other nodal points,  $\phi r$  takes the value 0, thus forms the polygonal line. Using  $\phi r$ , examine the polygonal line can be expressed in the form  $\sum_{r=1}^{M} C_r \phi_r$ , as well as in the form  $\int_0^1 \phi_r \phi_s dx = 0$ , 1(r = s, r = s).