Relativistic Electron Theory

Dirac Equation

Relativistic equation:

$$E^2 - p^2 c^2 = m^2 c^4$$

Define the condition for the equation above:

$$E \to i\hbar \frac{\partial}{\partial t} , \ p_x \to \frac{\hbar}{i} \frac{\partial}{\partial x}$$

Form a differential equation that makes the first-order differentiation in time:

$$\begin{split} i\hbar\frac{\partial\psi}{\partial t} &= H\psi \ , \ H = c\vec{\alpha}\cdot p + \beta mc^2 \ ; \ \vec{\alpha} \ , \ \beta = & \\ & \\ (i\hbar\frac{\partial}{\partial t} - c\vec{\alpha}\cdot p - \beta mc^2)\psi = 0 \end{split}$$
Usually Matrix

Operate on the conjugate operator from the left hand side:

$$(i\hbar\frac{\partial}{\partial t} + c\vec{\alpha} \cdot \boldsymbol{p} + \beta mc^2)(i\hbar\frac{\partial}{\partial t} - c\vec{\alpha} \cdot \boldsymbol{p} - \beta mc^2)\psi = 0$$

The following conditions for $\vec{\alpha}, \beta$ are then obtained through the comparison of the equation above and the equation $E^2 - p^2 c^2 - m^2 c^4 = 0$.

$$\begin{cases} \alpha_k \alpha_\ell + \alpha_\ell \alpha_k = 2\delta_{k\ell} \\ \\ \alpha_k \beta + \beta \alpha_k = 0 \\ \\ \beta^2 = 1 \end{cases}$$

To satisfy the condition, $\vec{\alpha}, \beta$ should be 4×4 matrix, and ψ is the vector in four-dimension. (from now on, $\vec{\sigma}$ represents Pauli matrix)

$$\beta = \begin{pmatrix} \mathbf{1} & \mathbf{0} \\ \mathbf{0} & -\mathbf{1} \end{pmatrix} , \ \vec{\alpha} = \begin{pmatrix} \mathbf{0} & \vec{\sigma} \\ \vec{\sigma} & \mathbf{0} \end{pmatrix}$$
$$\mathbf{1} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} , \ \mathbf{0} = \begin{pmatrix} 0 & 0 \\ 0 & 0 \end{pmatrix}$$

Dirac equation :
$$(i\hbar \frac{\partial}{\partial t} - c\vec{\alpha} \cdot \boldsymbol{p} - \beta mc^2)\psi = 0$$
, $\psi = \begin{pmatrix} \psi_1 \\ \psi_2 \\ \psi_3 \\ \psi_4 \end{pmatrix}$

$$P = \psi^{\dagger} \psi = \sum_{i} |\psi_{i}|^{2}$$

probability:

electric current density: $\mathbf{j} = \psi^{\dagger} \vec{\alpha} \psi$

Free electron:

$$\begin{split} \psi_j &= u_j e^{i \boldsymbol{k} \cdot \boldsymbol{r} - i w t} \\ E_{\pm} &= \pm \sqrt{c^2 p^2 + m^2 c^4} \\ E &= E_+ > 0: \quad u_+^{(1)} \sim \begin{pmatrix} 1 & 0 \\ 0 \\ c p_z / (E_+ + mc^2) \\ c (p_x + i p_y) / (E_+ + mc^2) \end{pmatrix} \sim \begin{pmatrix} 1 \\ 0 \\ v/c \\ v/c \end{pmatrix} \\ u_+^{(2)} &\sim \begin{pmatrix} 0 & 0 \\ 1 \\ c (p_x - i p_y) / (E_+ + mc^2) \\ -c p_z / (E_+ + mc^2) \\ -c p_z / (E_+ + mc^2) \end{pmatrix} \sim \begin{pmatrix} 0 \\ 1 \\ v/c \\ v/c \end{pmatrix} \\ E &= E_- < 0: \quad u_-^{(1)} \sim \begin{pmatrix} -c p_z / (-E_- + mc^2) \\ -c (p_x + i p_y) / (-E_- + mc^2) \\ 1 \\ 0 \end{pmatrix} \\ u_-^{(2)} &\sim \begin{pmatrix} -c (p_x - i p_y) / (-E_- + mc^2) \\ 0 \\ 1 \end{pmatrix} \end{split}$$

Free electron in electromagnetic field

$$\{i\hbar\frac{\partial}{\partial t} + e\phi - c\vec{\alpha} \cdot (\boldsymbol{p} + e\boldsymbol{A}) - \beta mc^2\}\psi = 0$$

$$\begin{pmatrix} \psi_3\\ \psi_4 \end{pmatrix} = \frac{c\vec{\sigma} \cdot (\boldsymbol{p} + e\boldsymbol{A})}{E + e\phi + mc^2} \begin{pmatrix} \psi_1\\ \psi_2 \end{pmatrix}$$

Where E > 0

$$\begin{split} E + e\phi + mc^2 &\approx 2mc^2 \ , \ \vec{\sigma} \cdot (\boldsymbol{p} + e\boldsymbol{A}) \ \vec{\sigma} \cdot (\boldsymbol{p} + e\boldsymbol{A}) = (\boldsymbol{p} + e\boldsymbol{A})^2 + e\hbar\vec{\sigma} \cdot \boldsymbol{H} \\ &\{\frac{1}{2m}(\boldsymbol{p} + e\boldsymbol{A})^2 + \frac{e\hbar}{2m}\vec{\sigma} \cdot \boldsymbol{H} - e\phi\} \begin{pmatrix} \psi_1 \\ \psi_2 \end{pmatrix} = E \begin{pmatrix} \psi_1 \\ \psi_2 \end{pmatrix} \end{split}$$

Central force field

$$\begin{split} H &= c\vec{\alpha} \cdot \boldsymbol{p} + \beta mc^2 + V(r) \\ E' &= E - mc^2 \\ [\frac{1}{2m(1 + \frac{E' - V}{2mc^2})}\boldsymbol{p}^2 + V - \frac{\hbar^2}{4m^2c^2}\frac{dV}{dr}\frac{\partial}{\partial r} + \frac{1}{2m^2c^2}\frac{1}{r}\frac{dV}{dr}\vec{\ell} \cdot \vec{s}] \binom{\psi_1}{\psi_2} = E'\binom{\psi_1}{\psi_2} \end{split}$$